The puzzling merging cluster Abell 1914: new insights from the kinematics of member galaxies

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ABSTRACT
We analyze the dynamical state of Abell 1914, a merging cluster hosting a radio halo, quite unusual for its structure. Our study considers spectroscopic data for 119 galaxies obtained with the Italian Telescopio Nazionale Galileo. We select 89 cluster members from spatial and velocity distributions. We also use photometry Canada-France-Hawaii Telescope archives. We compute the mean cluster redshift, $\langle z \rangle = 0.168$, and the velocity dispersion which shows a high value, $\sigma_V = 1210^{+125}_{-110}$ km s$^{-1}$. From the 2D analysis we find that Abell 1914 has a NE-SW elongated structure with two galaxy clumps, that mostly merge in the plane of the sky. Our best, but very uncertain estimate of the velocity dispersion of the main system is $\sigma_{V,\text{main}} \sim 1000$ km s$^{-1}$. We estimate a virial mass $M_{\text{sys}} = 1.4-2.6 \times 10^{15} \, h_{70}^{-1} \, M_\odot$ for the whole system. We study the merger through a simple two-body model and find that data are consistent with a bound, outgoing substructure observed just after the core crossing. By studying the 2D distribution of the red galaxies, photometrically selected, we show that Abell 1914 is contained in a rich large scale structure, with two close companion galaxy systems, known to be at $z \sim 0.17$. The system at SW supports the idea that the cluster is accreting groups from a filament aligned in the NE-SW direction, while that at NW suggests a second direction of the accretion (NW-SE). We conclude that Abell 1914 resembles the well-known nearby merging cluster Abell 754 for its particular observed phenomenology.

Key words: Galaxies: clusters: general. Galaxies: cluster: individual: Abell 1914.
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Optical information is an important way to investigate the dynamics of cluster mergers (Girardi & Biviano 2002). The spatial distribution and kinematics of galaxy members allow us to detect substructures and to analyze possible pre- and post-merging groups, and to distinguish between evolving mergers and remnants. Moreover, optical data are complementary to X-ray information because the ICM and galaxies react on different timescales during a collision. This is clearly shown in numerical simulations by Roettiger et al. (1997). Thus, for example, the importance of combining X-ray and optical data to study merging scenarios is shown by MUSIC (MULTi-Wavelength Sample of Interacting Clusters) project (Maurogordato et al. 2011).

In this context, we are now progressing on the DARC (Dynamical Analysis of Radio Clusters, see Girardi et al. 2010) project, which uses spectroscopic and photometric information of galaxy members to analyze the internal dynamics of clusters with diffuse radio emission.

We have carried out an intensive observational study focused on the cluster of galaxies Abell 1914 (hereafter A1914). A1914 is a rich cluster, X-ray luminous, hosting a hot ICM. It shows an Abell richness class $R = 2$ (Abell et al. 1989). $L_X(0.1-2.4 \text{ keV}) = 17.93 \times 10^{45} \, h^{-2} \text{ erg s}^{-1}$ (Ebeling et al. 1996) and $kT_X \sim \text{keV}$ (Baldi et al. 2007; Maughan et al. 2008). Following the Bautz-Morgan classification, A1914 is a type II structure (Abell et al. 1989), while it is a “L-type (”linear”) cluster in the Rood-Sastry morphological scheme (Struble & Rood 1987).

Dahle et al. (2002) study the mass and light distributions using weak-lensing techniques. They recover the elongated shape of this cluster in the NE-SW direction and find that the light distribution well follows the mass profile. The two brightest cluster galaxies trace the two highest peaks in the mass distribution, although in the reverse order, that is, the highest peak is close to the second brightest galaxy (Okabe & Umetsu 2008). On the other hand, Jones et al. (2005) analyze the galaxy distribution in the POSS digital. They find signs of dynamical activity with two distinct groups of galaxies with no single dominant galaxy.

Buote & Tsai (1996) develop the first analysis of the X-ray morphology of this cluster using ROSAT data. They find that A1914 is a relaxed structure, but Jones et al. (2005) show some evidence against this thesis, suggesting that this cluster is not so relaxed. They find no evidence for a cool core, an unusual high X-ray temperature, and notice that ROSAT data are very poorly fitted using a $\beta$ model. Then, using Chandra X-ray data, Govoni et al. (2004) show clear evidence of merger. In fact, they find a clear elongation of the X-ray surface brightness (along WNW-ESE, see Fig. 5d of Govoni et al. 2004). In the last years, using Chandra data, A1914 has been classified as a non relaxed cluster (Baldi et al. 2007; Maughan et al. 2008).

Concerning the radio emission, Komissarov & Gubanov (1994) first report evidence for a diffuse and extended radio source (see also Giovannini et al. 1999; Kempner & Sarazin 2001). Moreover, Bacchi et al. (2003), using VLA data, show the presence of a unpolarized halo. The halo covers a $7.4' \times 5.3'$ area with a power $P_{1.4 \text{GHz}} = 8.72 \times 10^{24} \, h_{70}^{-2} \text{ W Hz}^{-1}$. Govoni et al. (2004) point out that the diffuse radio emission is quite puzzling with a bright component elongated in the NW-SE direction and a more typical low-brightness halo in the cluster center (see Fig. 5d of Govoni et al. 2004). The bright radio region does not follows either the elongation of the X-ray surface brightness. This fact is quite unusual, because in the majority of clusters the elongated diffuse radio halo follows the direction of the merger (e.g., the “Bullet” Cluster 1E0657-56, Markevitch et al. 2002; but see Abell 523, Giovannini et al. 2011).

Despite several studies based on X-ray data, published redshift data are not enough to perform the detailed dynamical study of A1914. The work we present here is based on new spectroscopic data obtained with the Telescopio Nazionale Galileo (TNG). We also use photometric data from the Canada-France-Hawaii Telescope (CFHT) archive.

This paper is organized as follows. We present optical data, including redshifts and photometry information, in Sect. 2. We expose our results on the cluster structure in Sect. 3. The discussion on the dynamical state of A1914 and conclusions are presented in Sect. 4 and 5, respectively.

Unless otherwise stated, we present errors at the 68% confidence level (hereafter c.l.). Along this paper, we work using $H_0 = 70 \, \text{km s}^{-1} \text{Mpc}^{-1}$ and $h_{70} = H_0/(70 \, \text{km s}^{-1} \text{Mpc}^{-1})$ and a flat cosmology with $\Omega_M = 0.3$ and $\Omega_{\Lambda} = 0.7$. Within this cosmology, $\gamma$ corresponds to $\sim 172 \, h_{70}^{-1} \, \text{kpc}$ at the cluster redshift.

## 2 THE DATA SAMPLE

### 2.1 Spectroscopic data

We performed observations of A1914 using DOLORES multi-object spectrograph at the TNG telescope in March 2010. We used the LR-B grism, which provides a dispersion of 187 $\AA$/mm. DOLORES works with a 2048 × 2048 pixels E2V CCD. The pixel size is 13.5 $\mu$m. We retrieved a total of 4 MOS masks containing 146 slits. We exposed 3600 s for each mask.

Spectra were reduced using standard IRAF2 tasks. Radial velocities were computed by using the cross-correlation technique (Tonry & Davis 1979) with the IRAF/XCSAO task, as we have proceeded with other clusters already analyzed in the DARC project (for a detailed description, see

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1 see also http://adlibitum.oat.ts.astro.it/girardi/darc, the web site of the DARC project.

2 IRAF is distributed by the National Optical Astronomy Observatories, which are operated by the Association of Universities for Research in Astronomy, Inc., under cooperative agreement with the National Science Foundation.
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Figure 1. Multiwavelength image of the central region of A1914. The gray-scale image in background corresponds to the optical $r^{\text{Mag}}$ band (CFHT archive). Superimposed, with orange and yellow colors, we also show the smoothed X-ray image in the 0.3-7 keV energy range (Chandra archive). Contour levels represent the VLA radio image at 1.4 GHz (courtesy of F. Govoni; see Bacchi et al. 2003). Green circle and square mark the centers of the density peaks detected in our analysis of the galaxy distribution (see Sect. 3.6). Blue "X" marks the centroid of the X-ray surface brightness (Govoni et al. 2004). Labels indicate the positions of the two brightest cluster galaxies (BCG1 and BCG2; see Sect. 2.2 and Table 1). North is up and East is left.

In six cases (IDs. 78, 79, 83, 99, 100, 107 and 112; see Table 1), we considered the IRAF/EMSAO redshift (based on the wavelength of emission lines in the spectra) to get a realistic estimation. So, our catalog lists 113 galaxy redshifts in the field of A1914. We also considered six redshifts more from the SDSS archive (IDs. from 114 to 119; see Table 1).

The true intrinsic errors are larger than those formal errors given by the cross-correlation (e.g., Malumuth et al. 1992; Ellingson & Yee 1994; Quintana et al. 2000; Bardelli et al. 1994). To correct this effect, some galaxies were observed in more than one mask. This allows us to estimate the intrinsic errors in data of the same quality acquired with the same instrumentation. Our spectroscopic survey...
Figure 2. $r^\text{Mega}$-band image of Abell 1914 (CFHT archive). Circles and squares correspond to galaxy members and nonmembers, respectively (see Table 1). Labels and enhanced circles mark the positions of the two brightest cluster galaxies (BCG1 and BCG2; see Sect. 2.2 and Table 1). Black contours are the isodensity contours of the distribution of likely member galaxies (see Sect. 3.6 and Fig. 9-top panel). North is up and East is left.

provides duplicate estimations for 18 galaxies. So, following the method detailed in Barrena et al. (2009) for these 18 galaxies and using the weighted mean of the two measurement, we concluded that true intrinsic errors are larger than formal cross-correlation ones by a factor of two. For the redshifts estimated with the EMSAO task we considered the largest value between 100 km s$^{-1}$ and the formal error.

We also considered nine galaxies having redshift in the SDSS and lying in the same field spanned by our spectroscopic data. Three of these SDSS targets were also observed with the TNG/DOLORES. We find no systematic deviations between the SDSS and our redshifts. We add the remaining six galaxies to our TNG catalog. Finally, we obtain a spectroscopic catalog of 119 galaxies, with a median value of the $cz$ errors of 74 km s$^{-1}$.
2.2 Photometry and galaxy catalog
We also use public photometric data obtained with Megaprime/Megacam at the CFHT. In particular, we consider $g^\text{Mega}$ and $r^\text{Mega}$ band3 images retrieved from the CADC Megapipe archive (Gwyn 2009). These images cover an area of $1.05 \times 1.16 \, \text{deg}^2$ with a deepness of $g^\text{Mega} = 27.2$ and $r^\text{Mega} = 26.8$ limiting magnitudes (at the 5σ detection level). Megaprime photometry is 90% complete down to $g^\text{Mega} = 24.8$ and $r^\text{Mega} = 24.6$. We corrected $g^\text{Mega}$ and $r^\text{Mega}$ CFHT magnitudes for galactic extinction assuming extinction values obtained from SDSS (DR7) in the central cluster region.

Table 1 lists our velocity catalog and photometry (see also Fig. 2). We present an identification (ID) number in Col. 1 (galaxy members are listed in italic format); right ascension and declination (J2000) in Col. 2; CFHT $r^\text{Mega}$ magnitudes in Col. 3; and heliocentric radial velocities $v = c \zeta \Delta \lambda / c$ and errors $\Delta v$ in Col. 4 and 5, respectively.

Our spectroscopic sample is 80% (50%) complete down to $v = 18.3 \, \text{km s}^{-1}$ (within an elongated region of 60 arcmin$^2$ (corresponding to an area of $1.1 \times 1.6 \, \text{Mpc}$ at the redshift of the cluster) around the cluster center.

The brightest cluster galaxy is ID. 29 ($r^\text{Mega} = 15.87$, hereafter BCG1). It lies at the south-west and is non-dominant (in luminosity) in the cluster. In fact, there is a second brightest cluster galaxy at the north-east (ID. 20, $r^\text{Mega} = 16.39$, hereafter BCG2). The two galaxies are separated by $\sim 1.7 \, \text{Mpc}$ at the cluster distance.

3 ANALYSIS OF THE OPTICAL DATA
3.1 Cluster member selection
In order to select cluster members we followed a procedure with two steps. First, we run the 1D adaptive-kernel method (hereafter 1D-DEDICA, Pisani 1993 and 1996; see also Girardi et al. 1996; Fadda et al. 1996). We detected significant peaks (at >99% c.l.) in the velocity distribution. This procedure found A1914 as a peak at $z = 0.1675$, containing 100 (provisional) cluster candidates (in the range $45.921 \leq v \leq 58.839 \, \text{km s}^{-1}$, see Fig. 3). We also found 18 and 1 background and foreground galaxies, respectively.

Then, in a second step, we only consider the 100 likely candidate to run the “shifting gapper” method proposed by Fadda et al. (1996; see also, e.g., Girardi et al. 2011), which takes into account a combination of velocity and position of the galaxies. This method needs the definition of a cluster center, but the optical center of A1914 is not obvious due to the absence of a clear dominant galaxy and to the offset of the X-ray center (e.g., Maughan et al. 2008). So, we decided to assume as cluster center the position of BCG1 (see Table 1). The application of the “shifting gapper” rejected another eleven galaxies leading to a final sample of 89 cluster members (Fig. 4 – top panel).

In order to check the robustness in the galaxy member selection and estimate how the choice of cluster center could affect this selection, we executed the shifting gapper method, but now considering the BCG2 as cluster center. This procedure selected identical galaxy members. That is, in our case, the galaxy member selection, and so the dynamical analysis here exposed is not affected by the choice of the cluster center. So, we decided to consider the BCG1 as cluster center, as in agreement with Okabe & Umetsu (2008).

3.2 Global cluster properties
By using the biweight method (Beers et al. 1990, ROTAT software) with the 89 cluster members, we obtained a mean cluster redshift of $\langle z \rangle = 0.1678 \pm 0.0004$, i.e. $\langle v \rangle = (50.313 \pm 0.140) \, \text{km s}^{-1}$. In addition, by applying the same method and correcting for cosmological effects and
standard velocity errors (Danese et al. 1980), we obtained $\sigma_v = 1210^{+30}_{-110}$ km s$^{-1}$ (errors were estimated using the bootstrap technique). However, in order to check the robustness of this estimate, we study the variation of $\sigma_v$ with the distance to the cluster center (Fig. 4, bottom panel). The integral $\sigma_v$ profile is fl at, suggesting that the estimation of $\sigma_v$ is robust. Furthermore, when considering members within $0.1 h_70^{-1}$ Mpc from the BCGs (that is two groups of 7 and 8 galaxy members around BCG1 and BCG2, respectively), we measure similar mean velocities. We only find a modest difference, being $\langle v\rangle_{BCG2} < \langle v\rangle_{BCG1}$, which is in agreement with the finding that $v_{BCG2} < v_{BCG1}$ (see also section 3.5).

### 3.3 The small high velocity group

Figure 4 (top panel) shows that most of the interlopers have a very similar high velocity. We assign eight galaxies to a likely galaxy group (red squares in Fig. 4, top panel). Seven out of these eight galaxies lie in the southwest cluster region, thus reinforcing the idea that this is a real structure. For this high velocity galaxy group (hereafter HVG) we estimate $\langle v\rangle_{HVG} = (55557\pm89)$ km s$^{-1}$ and $\sigma_{v,HVG} = 221^{+55}_{-36}$ km s$^{-1}$.

### 3.4 Velocity distribution and 3D substructure

Deviations from Gaussianity in the velocity distribution are interpreted as an important sign that clusters present a complex dynamics (Ribeiro et al. 2011).

In order to check the Gaussianity in the velocity distribution, we used three profile estimators. These are the skewness, the kurtosis, and the scaled tail index STI (see Bird &
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Beers 1993). The STI finds evidence for non-Gaussianity at about 95%-99% c.l., suggesting a heavy tailed distribution (see Bird & Beers 1993 and their Table 2).

Furthermore, we investigated the presence of gaps in the velocity distribution. By using the weighted gap analysis presented by Beers et al. (1991; 1992; ROSTAT software), we detect one significant gap at the 97% c.l.. This gap divides A1914 into two groups, composed of 9 and 80 galaxies at low and high velocities, respectively. By applying the 2D Kolmogorov-Smirnov test (Fasano & Franceschini 1987) we see that the galaxies of the two groups present the same spatial distribution (see Fig. 5).

We also applied the 1D-Kaye’s mixture model (Ashman et al. 1994; see also, e.g., Boschin et al. 2012) – hereafter 1D-KMM – to search for bimodal partitions significantly fitting to the velocity distribution. The most likely solution (well under the 90% significance c.l.) indicates two groups of 15 and 74 galaxies, spatially not differing. Considering the KMM results, which take into account the group membership probability, we obtain that the two groups differ for about $\sim 140$ km s$^{-1}$ in the cluster rest-frame ($\langle v \rangle_{\text{KMM1D-LV}} = 50\,155$ and $\langle v \rangle_{\text{KMM1D-HV}} = 50\,319$ km s$^{-1}$) and the high velocity group has a much higher velocity dispersion ($\sigma_{V,\text{KMM1D-HV}} \sim 980$ km s$^{-1}$ vs. $\sigma_{V,\text{KMM1D-LV}} \sim 330$ km s$^{-1}$).

Correlations between spatial and velocity distributions of cluster galaxies usually indicate the presence of actual substructures. With this idea in mind, we used several techniques to reveal the structure of A1914 by combining positions and velocities. First, we searched for velocity gradients in the plane of the sky by performing multiple linear fits. The results of this test reveals no evidences of gradients. In addition, we performed a set of 3D tests: the classical $\Delta$ statistics (Dressler & Schechtman 1988), as well as its variation which considers separately mean velocity and velocity dispersion kinematical indicators (Girardi et al. 1997; Ferrari et al. 2003); the $\alpha$-test (West & Bothun 1990) and the $\epsilon$-test (Bird 1994) based on the projected mass predictions. None of the above mentioned tests yielded positive detection of substructures. Moreover, we found no substructure by applying the technique developed by Serna & Gerbal (1996), also named the ”Htree-method” (see also, Durret et al. 2010; Boschin et al. 2012).

3.5 2D galaxy distribution of the spectroscopic catalog

We applied the 2D adaptive-kernel technique (hereafter 2D-DEDICA) to the spatial distribution of member galaxies. This method found only one significant peak, lying close to BCG2 and elongated toward BCG1. Figure 7 shows that A1914 presents an elongated profile in the NE-SW direction, suggesting a bimodal structure. To further investigate this point we applied the 2D-KMM method using as seeds the cluster members contained within 0.1 $h^{-1}_{70}$ Mpc from BCG1.
3.6 2D galaxy distribution of the photometric catalogs

Our spectroscopic sample does not map the whole cluster field, besides of suffering for magnitude incompleteness. To overcome these restrictions we resorted to the photometric catalogs.

Using the CFHT photometry, we construct \((g'-r')\) color-magnitude relation (hereafter CMR, see Fig. 8), and select likely early-type members within the red sequence (RS) locus. In order to compute the RS, we apply a 2σ-clipping fitting procedure to the spectroscopic cluster galaxies. We obtained \(g^{\text{Mega}} \approx 1.341 - 0.021 \times r^{\text{Mega}}\) on 63 spectroscopic cluster members. With this method we selected as likely cluster members the objects within \(\pm 0.2\) mag with respect to the RS. Figure 8 shows as the selected magnitude intervals seem adequate to select RS galaxies. In this way, we only use good tracers of the cluster galaxy popula-
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4 DISCUSSION

We estimate a high value of the velocity dispersion, \( \sigma_V = 1210 \pm 110 \) km s\(^{-1}\). This result well agrees with a hot ICM showing a mean \( T_X \approx 9 \) keV (Baldi et al. 2007 and Maughan et al. 2008) when assuming energy equipartition between galaxies and gas energy per unit mass\(^4\), i.e. \( \beta_{\text{spec}} = 1 \), both suggesting a massive galaxy cluster. In the following sections we discuss our findings on the dynamical mass and cluster structure. Thus, based on these results, we propose a twobody model and time scale for the collision of substructures in Abell 1914.

4.1 Mass estimates

We computed the global virial quantities assuming the dynamical equilibrium (but see in the following) and in the framework of usual assumptions, i.e. cluster sphericity and coincidence in the galaxy-mass distributions. Following the method detailed in Girardi & Mezzetti (2001, see also Girardi et al. 1998) we obtained \( R_{200} \), an estimation of \( R_{200} \), and the mass within this radius. We assume a quasi-virialized region of \( R_{200} = 0.17 \times \sigma_V/H(z) \ Mpc \) (see Eq. 1 of Girardi & Mezzetti 2001 with the corresponding scaling of \( H(z) \) from Eq. 8 of Carlberg et al.

with \( r^{\text{Mega}} < 21 \), the two clusters closest in redshift (N.1 and N.2) are well detected, while the clusters N.3 and N.4 are better detected in the deeper magnitude ranges, thus confirming us that we are including more and more interlopers among our likely cluster members when considering fainter galaxies.

Table 2 lists information for the four highest, significant peaks in the galaxy distribution using the CFHT data: the estimated number of likely members, \( N_p \) (Col. 2); Equatorial coordinates of the substructure (Col. 3); the relative isodensity respect the highest peak, \( \rho_S \) (Col. 4); the \( \chi^2 \) for each clump (Col. 5). Galaxy clusters No.1 and No.4 are detected as significant peaks in the \( 21 \leq r^{\text{Mega}} < 22 \) sample, too, but having lower density.

Table 2. 2D substructure detected in the CFHT photometric data.

<table>
<thead>
<tr>
<th>2D – Subclump</th>
<th>( N_p )</th>
<th>( \alpha ), ( \delta ) (2000)</th>
<th>( \rho_S )</th>
<th>( \chi^2 )</th>
</tr>
</thead>
<tbody>
<tr>
<td>(( r = r^{\text{Mega}} ))</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>NE (( r &lt; 21 ))</td>
<td>167</td>
<td>26 : 02 : 2, 49 : 46</td>
<td>1.00</td>
<td>104</td>
</tr>
<tr>
<td>SW (( r &lt; 21 ))</td>
<td>176</td>
<td>25 : 52 : 0, 48 : 12</td>
<td>0.54</td>
<td>53</td>
</tr>
<tr>
<td>No.2 (( r &lt; 21 ))</td>
<td>204</td>
<td>25 : 02 : 2, 57 : 57</td>
<td>0.27</td>
<td>44</td>
</tr>
<tr>
<td>No.1 (( r &lt; 21 ))</td>
<td>99</td>
<td>24 : 49 : 4, 37 : 56</td>
<td>0.20</td>
<td>28</td>
</tr>
<tr>
<td>NE (( 21 \leq r &lt; 22 ))</td>
<td>90</td>
<td>26 : 06 : 7, 49 : 53</td>
<td>1.00</td>
<td>24</td>
</tr>
<tr>
<td>SW (( 21 \leq r &lt; 22 ))</td>
<td>96</td>
<td>25 : 53 : 0, 48 : 52</td>
<td>0.86</td>
<td>18</td>
</tr>
<tr>
<td>No.3 (( 21 \leq r &lt; 22 ))</td>
<td>46</td>
<td>26 : 30 : 8, 54 : 25</td>
<td>0.62</td>
<td>15</td>
</tr>
<tr>
<td>No.2 (( 21 \leq r &lt; 22 ))</td>
<td>80</td>
<td>25 : 30 : 4, 59 : 13</td>
<td>0.62</td>
<td>13</td>
</tr>
</tbody>
</table>

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Figure 9. Projected spatial distribution and isodensity contours (CFHT likely galaxy members with $r_{\text{Mega}} < 21$ (Top panel), $21 < r_{\text{Mega}} < 22$ (Middle panel), and $22 < r_{\text{Mega}} < 22.5$ (Bottom panel). Contours were estimated using the 2D-DEDICA method. "X" symbols mark the positions of BCG1 and BCG2. The numbers indicate the positions of the four galaxy clusters listed by NED within 20′ from the cluster center and having estimated or photometric redshift close to A1914 ($\Delta z < 0.06$), in order of increasing $\Delta z$. The circles indicate the cluster region within $2\, h_{70}^{-1}\, \text{Mpc}$ at the cluster distance, i.e., somewhat smaller than the virial radius ($2.2-2.7\, h_{70}^{-1}\, \text{Mpc}$, which is determined by the adopted model).

1997 for $R_{200}$). We estimate the mass using the equation $M = M_{\text{vir}} - SPT = 3\pi/2 \cdot \sigma_{V}^{2} R_{\text{vir}}/G - SPT$ (Eq. 3 of Girardi & Mezzetti 2001, with $R_{\text{vir}}$ derived as described in Eq. 13 of Girardi et al. (1998), with $A = R_{\text{vir}}$ as a close approach, and being the surface pressure term correction (SPT) a 20% of $M_{\text{vir}}$. Both $R_{\text{vir}}$ and $M$ were estimated considering the velocity dispersion with the usual scalings, where $R_{\text{vir}} \propto \sigma_{V}$ and $M(\sim R_{\text{vir}}) \propto \sigma_{V}^{3}$. We obtained $M(< R_{\text{vir}} = 2.7\, h_{70}^{-1}\, \text{Mpc}) = 2.6 \pm 0.8 \times 10^{15}\ h_{70}^{-1}\ M_{\odot}$.

On the other hand, it is commonly accepted that A1914 hosts a merging event, and so velocity dispersion and X-ray temperature could be enhanced (e.g., Ricker & Sarazin 2001; Schindler & Müller 1993). Our analysis fails to separate the cluster substructures in the velocity space. We have to assume the result derived from the 1D-KMM method (Sect. 3.4, i.e. the two best Gaussians obtained there, despite their low significance. According to these results, the secondary group presents a very small velocity dispersion and its mass can be neglected with respect to the main system. For the main system $\sigma_{V,\text{main}} \sim 980\ \text{km s}^{-1}$ which leads to $M(< R_{\text{vir}} = 2.2\, h_{70}^{-1}\, \text{Mpc}) = 1.4 \times 10^{15}\ h_{70}^{-1}\ M_{\odot}$. Hereafter, we consider as reliable the mass range $M_{\text{sys}} = 1.4 - 2.6 \times 10^{15}\ h_{70}^{-1}\ M_{\odot}$ for the whole A1914 system.

The above value of $\sigma_{V,\text{main}}$ is in agreement with $\sigma_{V,\text{SIS}} = 846 \pm 87$ estimated from the weak lensing analysis by Okabe & Umetsu (2008). For a punctual comparison with the projected mass computed by Okabe & Umetsu (2008) within $R=7\, \text{'},$ we project and rescale our mass estimate $M_{\text{sys}}$ assuming the cluster follows a NFW profile, taking a mass concentration parameter $c$ from Navarro et al. (1997) and correcting by the factor $1 + z$ (Bullock et al. 2001; Dolag et al. 2004, here $c \sim 1$). We obtain $M_{\text{sys,2D}}(< R = 1.2\, h_{70}^{-1}\, \text{Mpc}) = (1.1 - 2.0) \times 10^{15}\ h_{70}^{-1}\ M_{\odot}$, in agreement with that $M_{\text{2D}}(< R = 7\, \text{'}) = (4.10 \pm 1.55) \times 10^{14}\ h_{70}^{-1}\ M_{\odot}$ is a lower bound to the true enclosed mass (Okabe & Umetsu 2008). Instead, there is some tension between our value of $\sigma_{V,\text{main}}$ and the velocity dispersion estimate computed using redshifts from the on-going Hectospec Cluster Survey (Rines et al. 2010), $\sigma_{V,\text{Rines}} = 698^{+56}_{-38}\ \text{km s}^{-1}$ This explains the difference between our and their virial mass estimate $M_{\text{200, Rines}} \sim 6.0 \times 10^{14}\ h_{70}^{-1}\ M_{\odot}$, as obtained rescaling the value of $M_{\text{100}} (R_{\text{100}} \sim 1.3\, R_{200}$ and thus $M_{\text{200}} \sim 0.9 M_{\text{100}}$ for a NFW profile (Eke et al. 1996).

4.2 Cluster structure

Our 2D analyses confirm the existence of an important bimodal structure elongated in the NE-SW direction and find two significant peaks: the NE one closer to BCG2 and the SW one closer to BCG1, although not perfectly centered on the two BCGs.

The analysis of Govoni et al. (2004) shows that the X-ray peak is displaced with respect to the BCGs positions. They find that X-ray maximum is at South with respect to BCG2 (see their Fig. 5a). Similarly, we also find that the X-ray peak does not coincide with our peaks in the galaxy density (see our Fig. 1). The offset between the optical and X-ray peaks suggests a post-merger cluster but, as noted by Govoni et al. (2004), the X-ray features are not typical (e.g., simulations by Roettiger et al. 1998). In fact, their analysis shows the presence of a NE-SW arclike hot region crossing through the cluster center, while the X-ray emission is elon-
In this section, we present our efforts to unravel the dynamics of the merger between the two main subclusters in the SW-NE direction with a simple bimodal model, assuming that this collision causes (at least part of) the diffuse radio emission. Following the method detailed for other DARC clusters (see e.g., Abell 520 in Girardi et al. 2008 and Abell 2345 in Boschin et al. 2010), we apply the two-body model (Beers et al. 1982; Thompson 1982) to evaluate the timescales of the merger. This simple model assumes two point-mass bodies and a zero impact parameter. The model takes into account three parameters. These are the mass of the whole system, \( \mathcal{M}_{\text{sys}} \) (1.4–2.6 \( \times 10^{15} \) \( h_{100}^{-1} M_{\odot} \)), the relative line of sight velocity in the rest-frame, \( \Delta V \), and the projected linear distance between the two substructures, \( D \). As for the relative motion parameters, we consider our more reliable results, i.e. the estimate \( \Delta V = 140 \text{ km s}^{-1} \) obtained from our 1D analysis, and the estimate \( D \sim 0.4 h_{100}^{-1} \text{ Mpc} \) obtained from our 2D analysis. Due to the small radiative life of relativistic electrons and as a comparison with other radio halos clusters (e.g., Barrena et al. 2002; Girardi et al. 2008), we assume an elapsed time, \( t \), for the core crossing of few fractions of Gyr. We consider two different cases: \( t = 0.1 \) Gyr and \( t = 0.3 \) Gyr.

Figure 10 compares the model solutions as a function of \( \alpha \), where \( \alpha \) is the projection angle between the plane of the sky and the line connecting the centers of the two clumps, with the mass estimate of the system \( \mathcal{M}_{\text{sys}} \). At \( t = 0.1 \) Gyr, the solution is bound and outgoing (BO) with \( \alpha \sim 10^\circ \)–\( 25^\circ \), in agreement with the fact that we expect a merging axis mostly contained in the plane of the sky. At \( t = 0.3 \) Gyr, the model predicts unlikely angles \( \alpha > 60^\circ \). Even when considering as \( \mathcal{M}_{\text{sys}} \) the virial mass value by Rines et al. (2010, see our Sect. 4.1), the model with \( t \sim 0.1 \) Gyr should be preferred.

5 CONCLUSIONS

In conclusion, A1914 shows clear evidence of a recent cluster merger along the NE-SW direction and almost contained in the plane of the sky. The presence of an ongoing merger and the large mass are the main features of typical clusters with radio halos described in the literature. The large scale structure in the environment of A1914 suggests evidence of a second direction of cluster accretion, NW-SE. This merging axis is likely related to the bright feature of the diffuse radio emission. Thus, we argue that the unusual radio appearance of A1914 is due to the complexity of the merger. Indeed,
Figure 10. Two-body model applied to the NE and SW galaxy subclusters. The solutions are plotted as system mass vs. projection angle. Thick solid and thick dotted curves are bound and unbound solutions, respectively. Blue and red lines correspond to the case of $t = 0.1$ and $t = 0.3$ Gyr, respectively. $B_{1a}$ and $B_{1b}$ labels refer to bound and incoming, which denote the collapsing solutions (solid curve). On the other hand, the expanding solutions and unbound outgoing solutions (solid curve going on in the dotted curve, respectively) are labeled as BO and UO. Labels for the $t = 0.3$ case are skipped to keep the figure clear. Our mass estimate is enclosed by the horizontal solid lines, while the dashed line is the mass value obtained by Rines et al. (2010). The Newtonian criterion predicts a limit for bound solutions; this limit is shown as the thin dashed curve (above and below is the bound and unbound regimes, respectively).

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We point out that A1914 appearance resembles that of A754 where a complex merger scenario or a sloshing core seem to be the possible explanations. Only deeper X-ray data and many redshift measurements in a more extended cluster region could allow us to better understand the dynamics of A1914.

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